## PHY 555: Solid-state Physics I Homework #3

Due: 10/04/2024

## Conceptual

**1.** *5 points* What are the approximations made to go from the expression for the fully general tightbinding matrix elements and secular equation:

$$M_{ij\mathbf{k}} = \langle \Phi_{i\mathbf{k}} | H | \Phi_{j\mathbf{k}} \rangle$$
$$S_{ij\mathbf{k}} = \langle \Phi_{i\mathbf{k}} | \Phi_{j\mathbf{k}} \rangle$$
$$\det | M_{ij\mathbf{k}} - ES_{ij\mathbf{k}} | = 0$$

to the semi-empirical expressions:

$$M_{ij\mathbf{k}} = E_i \delta_{ij} + \sum_{\mathbf{t}_I} e^{i\mathbf{k}\cdot\mathbf{t}_I} \int \phi_i^*(\mathbf{r}) V_a(\mathbf{r} - \mathbf{t}_I) \phi_j(\mathbf{r} - \mathbf{t}_I) d^3 r$$
$$S_{ij\mathbf{k}} = \delta_{ij}$$
$$\det|M_{ij\mathbf{k}} - E\delta_{ij}| = 0$$

where  $\Phi_{i\mathbf{k}}$ 's are the Bloch sums,  $\phi_i$ 's are the atomic-like orbitals,  $\mathbf{t}_I$  runs over nearest neighbor lattice sites, and  $V_a(\mathbf{r} - \mathbf{t}_I)$  is the atomic-like potential centered on site  $\mathbf{t}_I$ .

## Analytical

**2.** 20 *points* Consider the empirical tight-binding treatment of *s* and *p* electrons on an FCC lattice. In this case, each lattice site has twelve nearest neighbors  $\mathbf{t}_l$ :  $(a/2)(0, \pm 1, \pm 1)$ ,  $(a/2)(\pm 1, 0, \pm 1)$ ,  $(a/2)(0, \pm 1, \pm 1)$ . The relevant interaction integrals are:

$$\int \phi_s^*(\mathbf{r}) V_a(\mathbf{r} - \mathbf{t}_I) \phi_s(\mathbf{r} - \mathbf{t}_I) d\mathbf{r} = V_{ss\sigma}$$

$$\int \phi_s^*(\mathbf{r}) V_a(\mathbf{r} - \mathbf{t}_I) \phi_{p_x}(\mathbf{r} - \mathbf{t}_I) d\mathbf{r} = l_x V_{sp\sigma}$$

$$\int \phi_{p_x}^*(\mathbf{r}) V_a(\mathbf{r} - \mathbf{t}_I) \phi_{p_x}(\mathbf{r} - \mathbf{t}_I) d\mathbf{r} = l_x^2 V_{pp\sigma} + (1 - l_x^2) V_{pp\pi}$$

$$\int \phi_{p_x}^*(\mathbf{r}) V_a(\mathbf{r} - \mathbf{t}_I) \phi_{p_y}(\mathbf{r} - \mathbf{t}_I) d\mathbf{r} = l_x l_y (V_{pp\sigma} - V_{pp\pi})$$

$$\int \phi_{p_x}^*(\mathbf{r}) V_a(\mathbf{r} - \mathbf{t}_I) \phi_{p_z}(\mathbf{r} - \mathbf{t}_I) d\mathbf{r} = l_x l_z (V_{pp\sigma} - V_{pp\pi})$$

where  $\phi_i$ 's are the atomic-like orbitals;  $V_{ss\sigma}$ ,  $V_{sp\sigma}$ ,  $V_{pp\sigma}$ , and  $V_{pp\pi}$  are the adjustable parameters; and  $\mathbf{l} = (l_x, l_y, l_z)$  is the unit vector in the direction of  $\mathbf{t}_I$ .

(a) Assume that the *s* orbitals do not interact with the *p* orbitals. Show that the dispersion of the *s*-orbital derived band is given by

$$E(\mathbf{k}) = E_s + 4V_{ss\sigma} \left[ \cos\left(\frac{k_x a}{2}\right) \cos\left(\frac{k_y a}{2}\right) + \cos\left(\frac{k_y a}{2}\right) \cos\left(\frac{k_z a}{2}\right) + \cos\left(\frac{k_z a}{2}\right) \cos\left(\frac{k_z a}{2}\right) \right]$$

where  $E_s$  are the onsite energies.

(**b**) For the *p* orbital manifold, show that the matrix elements are given by:

$$M_{p_x p_x \mathbf{k}} = E_p + 2\cos\left(\frac{k_x a}{2}\right) \left[\cos\left(\frac{k_y a}{2}\right) + \cos\left(\frac{k_z a}{2}\right)\right] (V_{pp\sigma} + V_{pp\pi}) + 4\cos\left(\frac{k_y a}{2}\right) \cos\left(\frac{k_z a}{2}\right) V_{pp\pi} M_{p_x p_y \mathbf{k}} = -2\sin\left(\frac{k_x a}{2}\right) \sin\left(\frac{k_y a}{2}\right) (V_{pp\sigma} - V_{pp\pi})$$

## Computational

**3.** 25 *points* Use the results of problem **2** to calculate the dispersion of the *s* and *p* bands in an FCC crystal along the path  $L \rightarrow \Gamma \rightarrow X \rightarrow K \rightarrow \Gamma$  (see Fig. 1 and Table 1 for high-symmetry **k** points/paths). Use the following parameters (all in Ha):  $V_{ss\sigma} = -0.5$ ,  $V_{pp\sigma} = 0.5$ ,  $V_{pp\pi} = -0.05$ ,  $E_s = 9$ ,  $E_p = 0$ . Use a = 10.67 Bohr for the lattice constant.



Figure 1: Brillouin zone and high-symmetry points/lines for the face-centered cubic Bravais lattice. ( $b_1$ ,  $b_2$ , and  $b_3$  correspond to the reciprocal lattice vectors of FCC, written below as  $\mathbf{g}_1$ ,  $\mathbf{g}_2$ , and  $\mathbf{g}_3$ .)

Table 1: High-symmetry **k** points of the face-centered cubic lattice.

-	$\times \mathbf{g}_1$	$\times \mathbf{g}_2$	$\times \mathbf{g}_3$
Γ	0	0	0
Κ	3/8	3/8	3/4
L	1/2	1/2	1/2
U	5/8	1/4	5/8
W	1/2	1/4	3/4
Х	1/2	0	1/2